

# **WHOLE -BRAIN AMIDE CEST IMAGING AT 3T WITH A STEADY-STATE RADIAL MRI ACQUISITION**

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# Abstract

To develop a steady-state saturation with radial readout chemical exchange saturation transfer (starCEST) for acquiring CEST images at 3 Tesla (T). The polynomial Lorentzian line-shape fitting approach was further developed for extracting amideCEST intensities at this field.

StarCEST MRI using periodically rotated overlapping parallel lines with enhanced reconstruction-based spatial sampling was implemented to acquire Z- spectra that are robust to brain motion. Multi-linear singular value decomposition postprocessing was applied to enhance the CEST SNR. The egg white phantom studies were performed at 3T to reveal the contributions to the 3.5 ppm CEST signal. Based on the phantom validation, the amideCEST peak was quantified using the polynomial Lorentzian line-shape fitting, which exploits the inverse relationship between Z-spectral intensity and the longitudinal relaxation rate in the rotating frame. The 3D turbo spin echo CEST was also performed to compare with the starCEST method.

The amideCEST peak showed a negligible peak B1 dependence between 1.2  $\mu$ T and 2.4  $\mu$ T. The amideCEST images acquired with starCEST showed much improved image quality, SNR, and motion robustness compared to the conventional 3D turbo spin echo CEST method with the same scan time. The amideCEST contrast extracted by the polynomial Lorentzian line-shape fitting method trended toward a stronger gray matter signal ( $1.32\% \pm 0.30\%$ ) than white matter ( $0.92\% \pm 0.08\%$ ;  $P = .02$ ,  $n = 5$ ). When calculating the magnetization transfer contrast and T1-

## ABSTRACT

corrected rotating frame relaxation rate maps, amideCEST again was not significantly different for white matter and gray matter.

Rapid multi-slice amideCEST mapping can be achieved by the starCEST method ( $< 5$  min) at 3T by combining with the polynomial Lorentzian line-shape fitting method.

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# Dedication

This thesis is dedicated to my parents, Keqing Yu and Jindong Sui, my friends and my cat, River, for their love and support.

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# Chapter 1 Introduction

CEST MRI is a molecular imaging technique that can detect proteins and metabolites through the water signal by making use of the exchangeable protons in these molecules.<sup>1-4</sup> Among endogenous CEST contrasts, amide proton transfer (or amideCEST) in mobile proteins has been intensively studied<sup>5-9</sup> due to its several favorable properties, including high concentration, large chemical shift, and optimal exchange rate ( $< 400$  Hz) for the NMR exchange regime at high and ultra-high field strengths (3 Tesla [T]-7T).<sup>5,6</sup> The successful acquisition of amideCEST signal requires sufficient SNR and image stability because both contrasts are only a few percent of the water signal and need to compete with stronger background signals from direct water saturation, and especially conventional semi-solid magnetization transfer contrast (MTC). Many efforts have been devoted to optimizing saturation parameters (e.g., saturation power and length) to increase CEST contrasts in tissues,<sup>10-13</sup> but improvements have been limited. Reducing the noise caused by oscillations in the Z-spectrum due to subject physiological and physical motions can also improve CEST SNR. This has been shown using motion-insensitive self-gated radial sampling CEST,<sup>14,15</sup> navigators to correct frequency fluctuations,<sup>16-19</sup> and/or some postprocessing methods by exploiting the redundancy of CEST images between different saturation offsets.<sup>20-23</sup>

Although radial CEST MRI has been successfully demonstrated for acquiring CEST and relayed nuclear Overhauser effect (rNOE) images on animals at high magnetic fields,<sup>14,15</sup> its application is



hindered by the long acquisition time on human scanners. In this work, we combined the fast radial acquisition method of periodically rotated overlapping parallel lines with enhanced reconstruction (PROPELLER),<sup>24,25</sup> with steady state saturation for amideCEST imaging on the human brain at 3T. PROPELLER is called MultiVane on the Philips system (Philips Healthcare, Best, The Netherlands). To further improve the steady-state saturation with radial readout chemical exchange saturation transfer (starCEST) SNR, multilinear singular value decomposition (MLSVD)<sup>26</sup> was applied. Recently, a radial sampling CEST method combining PROPELLER readout was proposed on rat brain.<sup>28</sup> However, due to the use of a conventional labeling scheme, that is, continuous-wave labeling followed by PROPELLER readout, this method was unable to reduce the total scan time.

The extraction of amideCEST signals from in vivo Z-spectra is another challenging issue for 3T MRI. The conventional CEST quantification method uses magnetization transfer ratio asymmetry analysis, that is, subtracting the labeling and control water saturation images acquired at the 2 symmetric offsets with respect to water resonance. However, amideCEST and aliphatic-rNOE resonances essentially are centered at opposite sides of the water frequency and, therefore, in part compensate each other when performing asymmetry analyses.<sup>5,28-31</sup> In order to distinguish between rNOE signals, we name them as aromatic-rNOE and aliphatic-rNOE in the current study. Furthermore, in white matter (WM) the MTC effect is known to be asymmetric with respect to water frequency<sup>33,34</sup> and also contributes. One popular strategy in the CEST field has been to acquire a full Z-spectrum with low saturation field strength B<sub>1</sub> and fit it by assuming a Lorentzian line-shape for each contributing signal, such as direct saturation (DS), guanidinium (Guan), amide, amine, and rNOE peaks.<sup>8,34-37</sup> This method can remove the semisolid tissue and DS components, and it works well in terms of extracting the rNOE signal. However, it remains challenging to extract the amideCEST signal due to the difficulty of removing interference from aromatic-rNOE

and other CEST peaks, such as the fast-exchanging amine and hydroxyl protons, to the CEST signal between 0 and 5 ppm. A similar strategy, named extrapolated semisolid magnetization transfer reference, was proposed for amideCEST/aliphatic-rNOE mapping with high saturation strengths; however, it still faces an issue similar to the Lorentzian line-shape fitting method when using low saturation strengths.<sup>38-40</sup> Recently, more evidence has suggested that the amide peak is discernible even at 3T MR fields when using low  $B_1$  strengths.<sup>41,42</sup> Hence, due to the smooth merging of background signals, the polynomial and Lorentzian line-shape fitting (PLOF) approach<sup>13,43,44</sup> that was proposed to extract and quantify CEST signals with discernible peaks should be well suited for amideCEST quantification imaging at 3T. In this work, we demonstrate that amideCEST imaging of the human brain at 3T is possible by combining the radial readout MRI with MLSVD in a clinically feasible acquisition time.

# Chapter 2 Method

## 2.1 Phantom preparation

Egg white is an appropriate model to demonstrate the pH sensitivity of the amide, Guan, and amine protons from mobile proteins. To demonstrate the possibility to detect pH-dependent amideCEST and rNOE signals with starCEST MRI, a set of egg white phantoms (100% liquid egg white) was prepared in a group of NMR tubes and titrated to pH values of 6, 6.5, 7, and 7.5 during slow stirring to avoid denaturation. To demonstrate the CEST contributions at 3.5 ppm, another set of egg white phantom was prepared by mixing egg white protein powder (NOW Sports Nutrition, Bloomingdale, IL) with pure water (0%  $D_2O$ ), 35%  $D_2O$  solution, and 70%  $D_2O$  solution, respectively. The egg white protein was fixed at 20% w/w with respect to the nondeuterated water in the solutions.

## 2.2 Human Studies

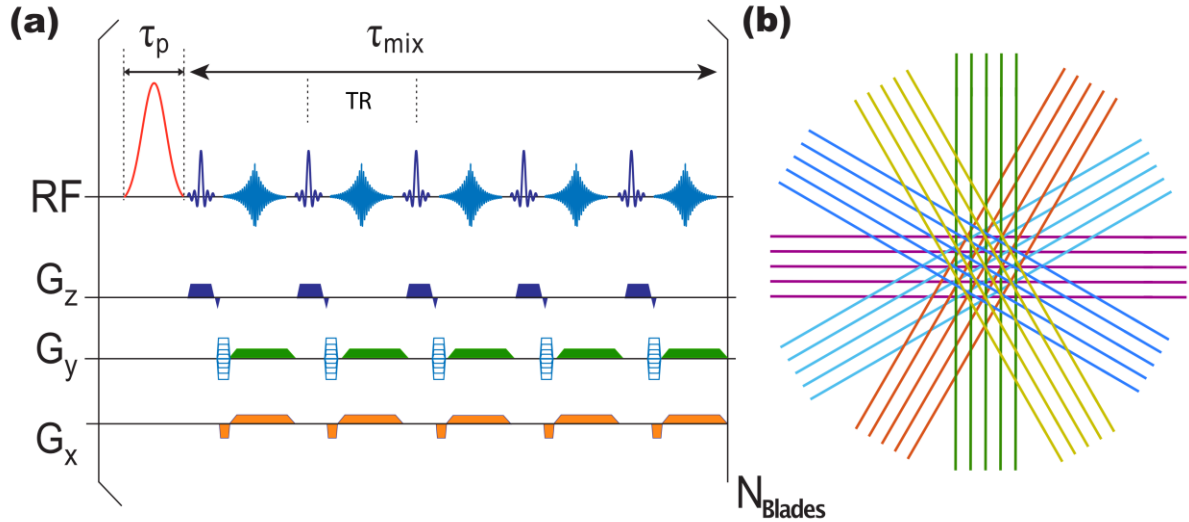
Human studies were approved by the Johns Hopkins Medicine Institutional Review Board and performed on 9 healthy subjects (23-46 years old), all of whom provided informed consent. All human scans were acquired on a 3T Philips Elition system (Philips Healthcare) using a quadrature transmit body coil and a 32-channel phased array receive head coil (Philips Healthcare). Two pads were placed on both side of the head close to the temporal lobes to minimize the head movement.

## 2.3 MRI experiments

The egg white phantom studies were performed to investigate the 3.5 ppm CEST contributions and to justify the proposed PLOF method for extracting the amideCEST at 3T. Considering several practical advantages over the human 3T, such as the high MRI image quality of the phantoms, the excellent  $B_0/B_1$  homogeneity, and the better temperature control with an air heater that is not available on the human scanners, the MRI experiments on the egg white phantoms were performed on a horizontal bore 3T Bruker Biospec system (Bruker, Ettlingen, Germany). During MRI experiments, the phantom was maintained at 37°C by an air heater. A 40-mm volume transceiver coil was used. The Z-spectrum was collected with an ultrashort echo time (UTE)-CEST sequence detailed previously.<sup>14,15,45,46</sup> The duration of the Gaussian saturation pulses was set to 50 ms with a peak  $B_1$  of 0.4-1.6  $\mu$ T unless specified, and the repetition time (TR) was 60 ms. The excitation pulse for the 2D-UTE radial readout was 10°. In order to make the peak  $B_1$  comparable to the  $B_1$  values used in the conventional continuous wave CEST, the equivalent block pulse  $B_1$  values ( $B_{1ave}$ ) were also reported by  $B_{1ave} = A_{area} \times B_1$ , where  $A_{area}$  is the area ratio between 1 sinc-Gauss ( $A_{area} = 0.54$ ) or Gaussian ( $A_{area} = 0.42$ ) pulse and 1 block pulse with the same peak  $B_1$ . Here, the averaged  $B_{1ave}$  value was only for each saturation pulse (length  $\tau_p$ ) and was not the averaged value for the total period  $\tau_p + \tau_{mix}$  that determines the steady state. The effective echo time (TE) was 0.3 ms. In the current study, single-slice CEST images were collected with a matrix size of  $96 \times 96$  and a slice thickness of 5 mm. The total imaging time for each offset was 12 s. The saturation offset was swept from -12 to 12 ppm using a 0.2-ppm increment between -8 and 8 ppm and a 0.5-ppm increment for other regions.

Due to the limited acquisition speed of the UTE-CEST-based radial readout method,<sup>46,47</sup> we applied the PROPELLER-based method for the acquisition of multi-slice CEST on the 3T clinical

scanner. Figure 1a depicts the starCEST sequence based on the MultiVane Philips version of this sequence. Notice that we plotted 5 blade lines and echoes for illustration in Figure 1, but we acquired 12 blade lines in practice. Every TR contains a sinc-Gauss saturation pulse (50 ms) followed by a train of 12 gradient echoes for detection (TR/ TE/flip angle = 4.42 ms/2.08 ms/10°), which generates the 12 blade lines in 1 blade. Each blade was rotated and shifted following the motion correction method proposed in the literature,<sup>25</sup> and then the images were reconstructed by gridding them to a Cartesian data set. With fixed TR, the number of lines per blade determines the mixing time between saturation pulses, following  $\tau_{mix} = TR \times (\#linesperblade)$ . Previous pulsed CEST optimization has already demonstrated that 1 pulse width larger than 40 ms would achieve enough frequency selectivity at 3T.<sup>49</sup> Hence, the pulse width of the sinc-gauss saturation pulses was set to 50 ms. For the human brain study, images were acquired at  $5 \times 3 \times 3 \text{ mm}^3$  resolution



**Figure 1: The illustration of starCEST**

(a) Timing diagram of the starCEST sequence based on PROPELLER.  $\tau_p$  is the duration of the saturation pulses;  $\tau_{mix}$  is the mixing time for saturation transfer between successive saturation pulses; TR is the repetition time between 2 k-line acquisitions, determined by the number of phases encodes per k-line. (b) Illustration of the sampling of k-space in starCEST using a set of radially directed blades. In this example, for demonstration purposes only, each blade is composed of 5 parallel phase-encoded lines that are collected using gradient echoes. However, 12 phase-encoded lines were acquired per blade in the current paper. starCEST, steady-state saturation with radial readout Chemical Exchange Saturation Transfer, PROPELLER, periodically rotated overlapping parallel lines with enhanced reconstruction.

using 15 slices and a FOV of  $220 \times 220 \text{ mm}^2$ . The image matrix was resized to  $96 \times 96$  during reconstruction. The parallel imaging SENSE factor was 2 in the anteroposterior direction for the lines in each blade. The oversampling rate was 150% for MultiVane (Philips Healthcare), and 12 blade lines in 1 blade were collected unless specified. Hence, 9 blades were collected for each slice. The time for the multi-slice acquisition was 7.6 s per irradiation frequency. Furthermore,<sup>37</sup> saturation frequency offsets from 1 to 7.8 ppm with 0.2-ppm increment were acquired for amideCEST imaging with a total acquisition time of 4min40s. The nonsaturated reference signals were acquired using the same sequence with a saturation offset of 200 ppm. The second-order pencil beam-volume shimming method was applied on the image slices before all the CEST experiments.

The conventional 3D turbo spin-echo (TSE) CEST method<sup>38,49,50</sup>, was performed on 2 subjects together with the starCEST method to compare the image quality, SNR, and motion robustness. The CEST preparation module was identical to the saturation scheme in the starCEST, that is, 50 ms sinc-Gauss saturation pulses with a peak  $B_1$  of 1.4  $\mu\text{T}$  and 53 ms mixing times. Thirty-two pulses and mixing times were implemented, which resulted in a total saturation time of 3.3 s and were considered as a steady-state saturation. The matrix and geometry sizes were identical to those used in the starCEST. The TSE factor was 174 with a 120-degree refocused angle.  $\text{TR} = 7.2 \text{ s}$ ,  $\text{TE} = 12 \text{ ms}$ , and the compressive sensing factor = 4 were used. The offsets and the total experimental time were identical to those of starCEST. In order to demonstrate the impact of brain motions on the amideCEST maps, starCEST and TSE experiments were performed sequentially by removing the 2 pads that was used for fixing the brain.

$T_1$  maps were acquired with a look-locker sequence with geometry identical to the CEST scan and were used for the CEST quantification in the PLOF method.<sup>51</sup> Furthermore, 38 images were

acquired with TR = 157 ms (TI = 35-5870 ms) and a readout flip angle of 4°. Pixel-by-pixel fitting was performed to a 3-parameter model  $\left(A - Be^{-\frac{TI}{T_1^*}}\right)$  to obtain  $T_1$  (as  $T_1 = (B/A - 1)T_1^*$ ).<sup>52</sup>  $B_0$  maps were obtained using a dual-echo sequence with TR = 10 ms, TE = 4.6 ms, and a flip angle of 30°.  $B_1$  maps were obtained using the dual refocusing echo acquisition mode technique with a stimulated echo acquisition mode (STEAM) at a flip angle of 60°.<sup>53</sup>

## 2.4 Data analysis

All MRI images were processed using custom-written MATLAB scripts (MathWorks, Natick, MA; www.mathworks.com, version 2020b). The amideCEST signal from both starCEST and 3D TSE CEST method was extracted using the PLOF method, which has been described previously.<sup>13,43,44</sup> Briefly, the steady-state normalized Z-spectral intensity  $Z^{SS}$  has an inverse relationship to the longitudinal relaxation rate in the rotating frame,  $R_{1\rho}$ , that is, the rotating-frame relaxation spectrum (R-spectrum)<sup>53-56</sup>:

$$Z^{SS}(R_{1\rho}) = \frac{\cos^2 \theta R_1}{R_{1\rho}} . \quad (1)$$

$R_1 = 1/T_1$  is the longitudinal relaxation rate of water, and  $\theta$  is the tilt angle of the effective magnetization with respect to the Z-axis.  $R_{1\rho}$  is the water relaxation rate during saturation, which includes contributions from the effective water relaxation rate  $R_{eff}$ , the rotating frame rate from the CEST signal of interest ( $R_{exch}$ ), and a term  $R_{back}$  that accounts for the rotating frame rates of all other magnetization/exchange transfer processes in tissue<sup>55</sup>:

$$R_{1\rho} = R_{eff} + R_{exch} + R_{back} , \quad (2)$$

where  $R_{eff}$  is the longitudinal relaxation rate of water in the rotating frame without additional solution components. The observed CEST signal  $\Delta Z$  is extracted by the following:

$$\Delta Z = Z(R_{eff} + R_{back}) - Z(R_{eff} + R_{exch} + R_{back}). \quad (3)$$

The amideCEST peak ( $R$ ) and broad background ( $R_{back}$ ) in the R-spectrum can be represented by a Lorentzian function and a mixed polynomial and Lorentzian function, respectively:

$$R = R_{amide} \frac{(w/2)^2}{(w/2)^2 + (\Delta\omega - \Delta\omega_{amide})^2} \quad (4)$$

$$R_{back} = \frac{C_0(C_1/2)^2}{(C_1/2)^2 + \Delta\omega^2} + C_2 + C_3 \cdot \Delta\omega, \quad (5)$$

where  $w$  is the peak full width at half maximum of the Lorentzian line-shape.  $\Delta\omega$  is the offset relative to the water proton resonance.  $R_{amide}$  is the intensity of the amideCEST peak in the R-spectrum.  $\Delta\omega_{amide}$  is the chemical shift offset of the CEST peak relative to the water signal, that is, 3.5 ppm; The background function was modified from the original polynomial function<sup>13,43,44</sup> to improve the PLOF background fitting at 3T. The first Lorentzian function in Equation 5 was included to account for the water DS, whereas the polynomial function ( $C_2 + C_3 \cdot \Delta\omega$ ) was implemented to fit the MTC and other exchanging protons such as amine/hydroxyl CEST. In the current study, we found that a linear function can perfectly fit the combined effect with MTC and other CEST between 1 and 7.6 ppm. For a wider fitting range, a higher order approximation can be used. The Z-spectral range of [2.5, 4.5] (ppm) was excluded from the background fitting (fitting range = 1.0-7.6 ppm) for amideCEST. When performing the PLOF calculation on the brain, the measured  $B_0$  map was used to shift the Z-spectrum for field homogeneity correction in each pixel. The MLSVD method was applied to enhance the SNR for amideCEST maps.<sup>26</sup> The truncation numbers were 48, 48, and 10, and they were used for the 2 image dimensions and Z-spectral dimension, respectively. The SNR of the images were calculated by the mean values in the signal

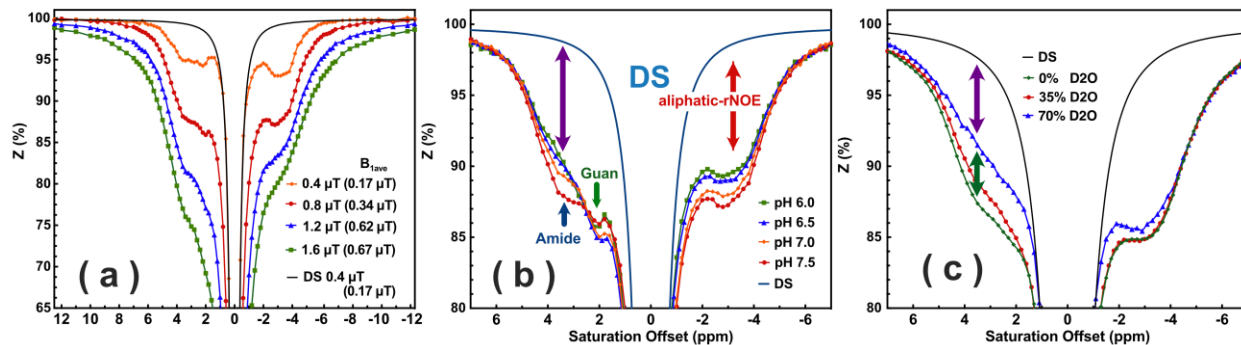


regions divided by the standard deviation (SD) of the noise regions at the 4 image corners. The intersubject reproducibility was quantified by coefficient of variance (coefficient of variance =  $SD/mean \times 100\%$ ). The intraclass correlation coefficient was used to measure the intrasubject agreement between 2 different scans.<sup>57</sup> The coefficient of variance and intraclass correlation coefficient for the amideCEST ( $\Delta Z_{amide}$ ) were calculated based on the averaged values in the whole brain gray matter (GM) and WM regions.

# Chapter 3 Results

## 3.1 Phantom

We first studied the spectral contributions and pH dependence of the egg white Z-spectrum using UTE-CEST as a function of pH values. Figure 2 displays the Z-spectra of egg white as a function of saturation strength and pH. At low peak  $B_1$  values ( $< 0.8 \mu\text{T}$  or  $B_{1\text{ave}} < 0.43 \mu\text{T}$ ), the amide, Guan, and aliphatic peaks show structure, whereas those peaks merge with the DS lineshape at higher saturation field strengths ( $> 1 \mu\text{T}$ ). The results also demonstrate the frequency range of



**Figure 2: The Z-spectrum of egg white**

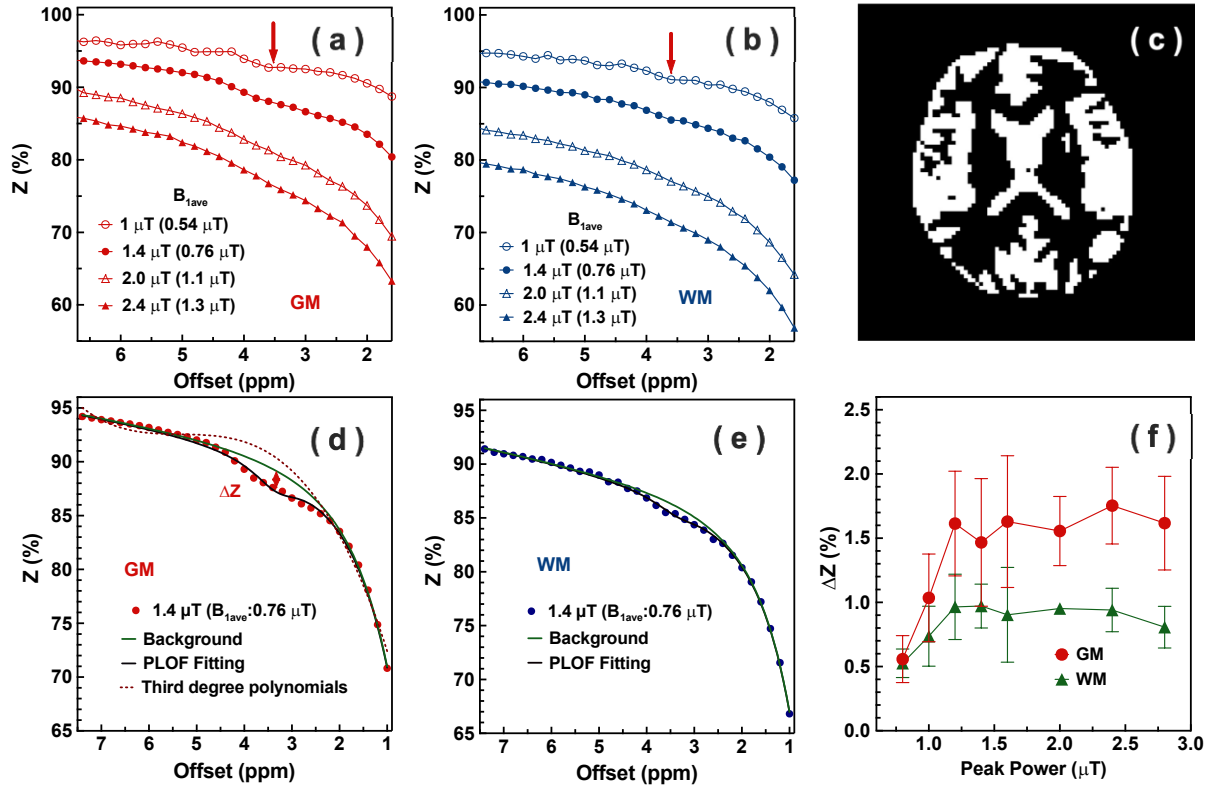
(a) Z-spectrum of egg white (pH 7.0) as a function of peak  $B_1$ . A simulated DS spectrum for peak  $B_1 = 0.4 \mu\text{T}$  ( $B_{1\text{ave}} < 0.17 \mu\text{T}$ ) is also shown (black curve). (b) Z-spectrum of egg white as a function of pH with peak  $B_1 = 0.8 \mu\text{T}$  ( $B_{1\text{ave}} < 0.34 \mu\text{T}$ ). The exchangeable protons downfield, amide and Guan, and aliphatic-NOE signals are indicated. The simulated DS spectrum with peak  $B_1 = 0.8 \mu\text{T}$  ( $B_{1\text{ave}} < 0.34 \mu\text{T}$ ) is also plotted. (c) The Z-spectra of egg white solution with 0% (green), 35% (red), and 70% (blue)  $D_2O$  acquired with UTECEST with peak  $B_1 = 0.8 \mu\text{T}$  ( $B_{1\text{ave}} < 0.34 \mu\text{T}$ ). A residual signal is still observable in the exchangeable proton range for 70%  $D_2O$  (purple arrow). The changes in the Z-spectral broad background at 3.5 ppm between 0% and 70%  $D_2O$  (green arrow) are attributed to CEST effects of amine, and other fast exchangeable protons (e.g., hydroxyl in myo-inositol) that are partially merged with the water signal due to being in the intermediate to fast exchange regime at 3T. The simulated DS spectrum for egg white in 0%  $D_2O$  is also plotted. DS, direct saturation; Guan, guanidinium; NOE, nuclear Overhauser effect; T, Tesla

the protein Z-spectrum, which widens as a function of saturation field strength<sup>42</sup> due to the widening of the CEST peaks and larger bandwidths of the pulses. At 3T, the protein solution Z-spectrum range is approximately  $\pm 9$  ppm for peak  $B_1 = 0.8 \mu\text{T}$  ( $B_{1\text{ave}} < 0.43 \mu\text{T}$ ). For the amideCEST peak around 3.5 ppm, decreased effect size was found for lower pH values (Figure 2b); however, relative to the estimated DS a strong residual saturation transfer background signal remained over a broad frequency range even at pH 6 (purple arrow in Figure 2b). As expected for exchange-relayed NOEs and shown previously in phantoms,<sup>58,59</sup> the aliphatic rNOE at the opposite frequency range also demonstrated pH dependence. The Guan effect size first increased when lowering pH from 7.5 to 7, and then it stabilized and decreased with further reduction of pH, which was consistent with previous observations.<sup>60</sup> The contributions of CEST and rNOE signals in protein were examined by varying the hydrogen/deuterium ratio of water in egg white solutions, and the results at a peak  $B_1$  of  $0.8 \mu\text{T}$  ( $B_{1\text{ave}} < 0.43 \mu\text{T}$ ) are plotted in Figure 2c. As expected, the aliphatic-rNOE showed only a small reduction, whereas the amide peak and part of the broad background in the exchangeable proton range were clearly reduced for the 35%  $D_2O$  solution and even more for the 70%  $D_2O$  solution. For 70%  $D_2O$ , the CEST effect between 0 and 6 ppm decreased dramatically, but still a residual signal (about 56% in the 20% egg white solutions) was observable (purple arrow in Figure 2c).

## 3.2 Human brain

Examples of starCEST Z-spectra between 1.6 and 6.7 ppm for the human brain recorded as a function of saturation strength are depicted in Figure 3a,b for GM and WM, respectively. Spectra are from regions of interest drawn on the brain in the mid-axial slice (8th slice) for which the  $T_1$ -based mask for separating WM and GM is shown in Figure 3c. A discernable amide peak at 3.5 ppm can be observed for peak  $B_1$  strengths below  $2 \mu\text{T}$  ( $B_{1\text{ave}} < 1.1 \mu\text{T}$ ) for both GM and WM

and the peak merges with the background above 2  $\mu\text{T}$ . The Z-spectra acquired with a peak  $B_1$  of 1.4  $\mu\text{T}$  from Figure 3a,b are shown in Figure 3d,e and demonstrate obvious differences in the amide peak for GM and WM; that is, the amide peak in GM is more pronounced than that in WM. For example, the amide peak in GM is more pronounced than that in WM. The MTC background of WM is stronger  $\mu\text{T}$  than in GM; that is, GM absolute Z-values are higher than those of WM, which

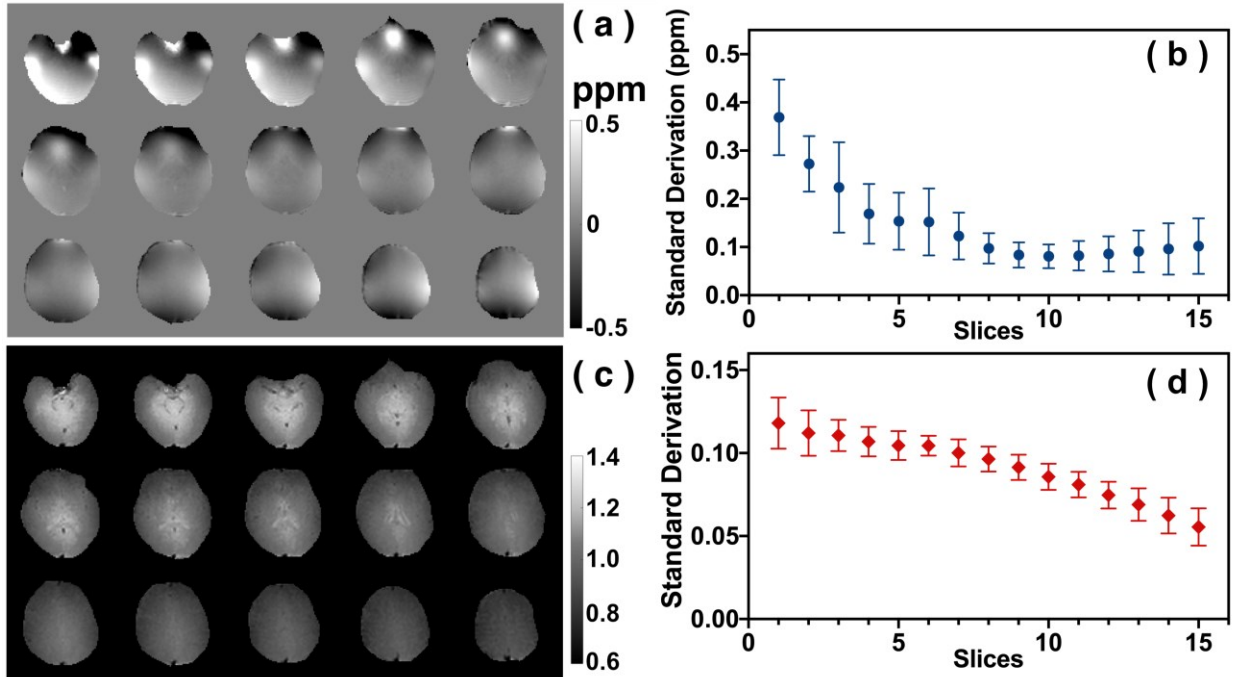


**Figure 3: Typical Z-spectrum of human brain**

Typical Z-spectrum of human brain acquired with starCEST on a 3T clinical scanner as a function of peak  $B_1$  for gray matter (GM) (a) and white matter (WM) (b), respectively. The regions of interest (GM and WM) were extracted from a mid-axial slice (8th slice) using a mask (c) based on the T1 map with a cut-off value of 0.8 s. (d) A GM Z-spectrum of the mid-axial slice recorded with peak  $B_1 = 1.4 \mu\text{T}$  ( $B_{1\text{ave}} < 0.76 \mu\text{T}$ ) and the PLOF fitting results with Equations 1-5. The method to extract  $\Delta Z_{\text{amide}}$  signal with Equations 1-5 is indicated using a red arrow. The PLOF background fitted with the conventional 3rd degree polynomial function<sup>13,43,44</sup> was plotted for comparison (red dashed line). The amide proton range and the whole fitting range were identical for the 2 background functions (Equation 5 and the 3rd degree polynomial function). (e) A WM Z-spectrum of the mid-axial slice recorded with peak  $B_1 = 1.4 \mu\text{T}$  ( $B_{1\text{ave}} < 0.76 \mu\text{T}$ ) together with the PLOF fitting results. (f)  $B_1$  dependence of the amideCEST signal for the WM and GM in the mid-axial slices ( $n=4$ ). Error bars reflect the SD over different subjects. The signal was extracted with is the PLOF method illustrated in (d) and (e). GM, gray matter; PLOF, polynomial and Lorentzian line-shape fitting; WM, white matter

is consistent with conventional MTC studies. In order to illustrate the process of extracting the amideCEST signal with the PLOF method, the PLOF fitting curves for the background  $Z_{back}$  with the improved background function and the full Z-spectrum are plotted in the same figures (Figure 3d,e). The background can be well fitted with 4 parameters in Equation 5 for the whole 1-7.5 ppm, which allowed us fully utilizing all acquired offsets. The resulting  $B_1$ -dependent amideCEST signal is shown in Figure 3f and did not show clear peak  $B_1$  dependence between 1.2 and 2.4  $\mu\text{T}$  ( $B_{1ave}$  between 0.65 and 1.3  $\mu\text{T}$ ) for both WM and GM. In order to make the sequence robust to the  $B_1$  variation in the brain, the peak  $B_1$  was therefore set to 1.4  $\mu\text{T}$  ( $B_{1ave} < 0.76 \mu\text{T}$ ) in the following amideCEST studies.

Figure 4 depicts  $B_0$  and  $B_1$  maps of the brain of a typical subject. The  $B_0$  and  $B_1$  SD for each slice are also plotted in Figure 4b,d, respectively. The  $B_0$  inhomogeneity over the whole brain was

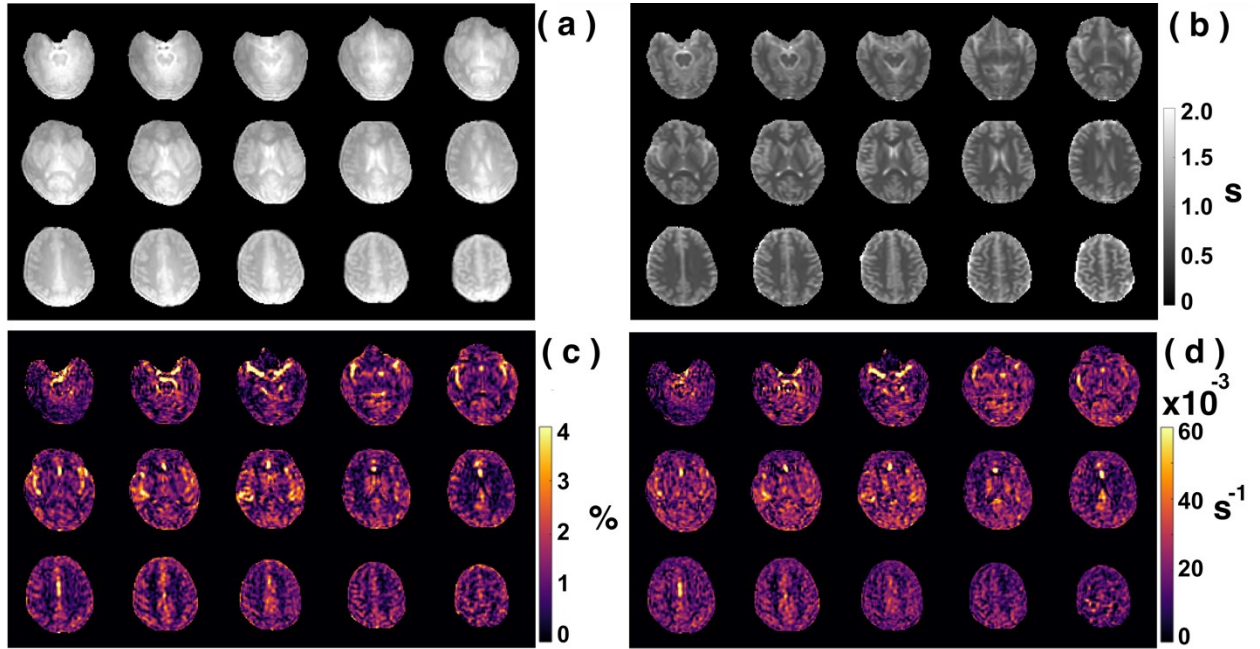


**Figure 4: Multi-slice  $B_0$  and  $B_1$  maps**

(a) Typical multi-slice  $B_0$  maps acquired with dual-echo sequence on the human brain at 3T. (b) The averaged SD ( $n=3$ ) for each slice of the  $B_0$  maps. (c) Typical multi-slice  $B_1$  ratio maps obtained with DREAM method on the same subject. (d) The averaged SD ( $n=3$ ) for each slice of the  $B_1$  maps. DREAM, dual refocusing echo acquisition mode; Ramide, rotating frame relaxation rate of amide protons.

in the range of  $-0.7$ - $0.7$  ppm. The major inhomogeneity areas ( $STD > 0.2$ ) appeared at the beginning of the slice package (slices 1-3). The  $B_1$  map was quite homogeneous across the brain ( $STD < 0.11$ ), except for the typical  $B_1$  hyperintensity in the midbrain region (slices 1-3). Figure 5 depicts starCEST,  $T_1$ ,  $\Delta Z_{amide}$ , and  $R_{amide}$  maps of the brain of a typical subject. The  $\Delta Z_{amide}$ , and  $R_{amide}$  maps were calculated with the PLOF method by including the  $B_0$  map and the measured  $T_1$  maps. Discernable WM and GM amide CEST effects can be observed in the  $\Delta Z_{amide}$  maps, and the GM amideCEST ( $\Delta Z_{amide}$ ) is higher than that of WM, consistent with the observation from typical WM and GM Z-spectra (Figure 3d,e). Correction for the MTC/DS background and  $T_1$  impacts the rotating frame relaxation method (Equations 1-5), and the WM and GM contrast in the  $R_{amide}$  maps diminished significantly.

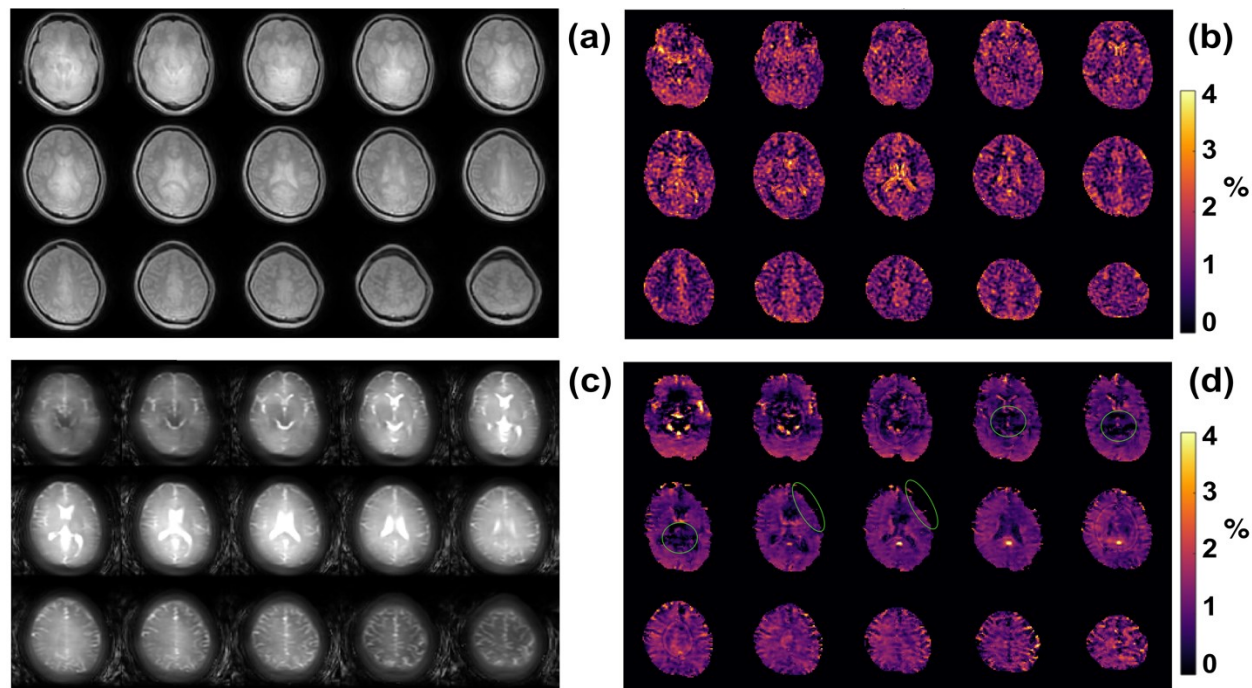
The whole brain images acquired with the starCEST and 3D TSE CEST are plotted in Figure 6 together with the corresponding amideCEST maps. The starCEST images show much better image



**Figure 5: starCEST,  $T_1$ ,  $\Delta Z_{amide}$ , and  $R_{amide}$  maps of the brain**

(a) starCEST images; (b)  $T_1$  maps by look-locker sequence. The amideCEST maps ( $\Delta Z_{amide}$ ) and the corresponding apparent relaxation rate ( $R_{amide}$ ) maps extracted with the PLOF method by including the  $B_0$  and  $T_1$  maps are shown in (c) and (d), respectively.

quality than the TSE CEST images as judged from the clear contrast between WM/GM (Figure 6a). In contrast, the images obtained by the TSE method are blurry due to the long TSE acquisition time (2.1s), and it is challenging to differentiate WM and GM regions. Strong CSF signal was observed in the TSE images (Figure 6c). The starCEST SNR values in the center slice were much higher than those of TSE CEST images (starCEST GM vs. TSE GM:  $96 \pm 19$  vs.  $54 \pm 16$ ; starCEST WM vs. TSE WM:  $34 \pm 6$  vs.  $18 \pm 6$ ). The amideCEST maps obtained by starCEST (Figure 6b) did not show any obvious motion introduced artifacts in both GM and WM regions, but hyperintensities were present in some of the CSF regions. For the amideCEST maps obtained by the TSE CEST method, serious hypo intensities were observed in the most of slices, tentatively attributed to the brain motions (Figure 6d).

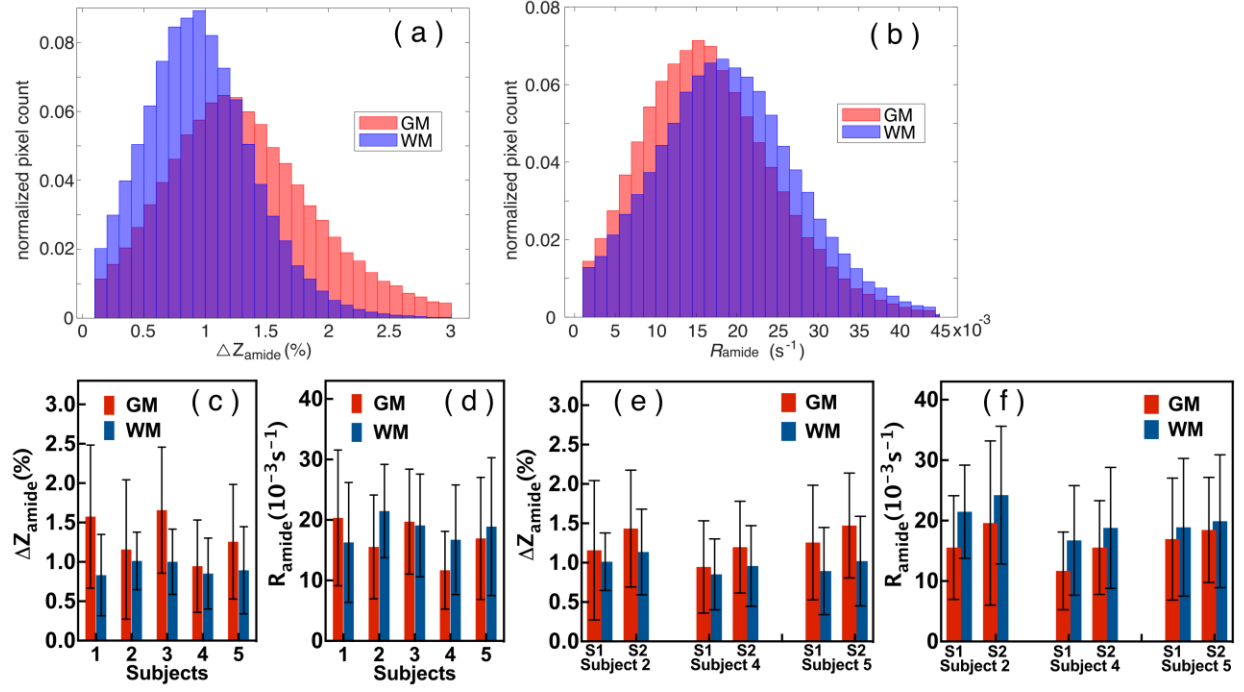


**Figure 6: Whole brain images and amideCEST maps acquired with the starCEST and 3D TSE CEST**

Comparison of the human brain images acquired at 3T with the multi-slice starCEST (a) and 3D TSE CEST sequence (c). The amideCEST maps extracted with the PLOF method from the multi-slice starCEST (b) and 3D TSE CEST (d) images. Hypo intensity regions in the amideCEST maps due to the brain motions are indicated with green circles. TSE, turbo spin-echo

In order to quantify the amideCEST values, Figure 7 shows histograms of the amideCEST difference and rotating frame relaxation values in WM and GM for all subjects ( $n = 5$ ). The regions of interest for GM and WM were extracted from the T1 maps with a cutoff value of 0.8 s. The mean and SD of the amideCEST signals were  $1.32\% \pm 0.30\%$ (GM) and  $0.92\% \pm 0.08\%$  (WM) ( $P=.02$ ,  $n=5$ ) for  $\Delta Z_{amide}$ , respectively (Figure 7a). In the Ramide images, the averages also showed no significant difference (GM:  $17 \pm 4 \times 10^{-3} \text{ s}^{-1}$ ; WM:  $18 \pm 2 \times 10^{-3} \text{ s}^{-1}$ ;  $P=.4$ ,  $n=5$ ) (Figure 7b), consistent with the low contrast in the Ramide maps (Figure 5d). Reproducibility was investigated for the same protocol across different subjects ( $n = 5$ ) as well as the same subject ( $n = 3$ ) at 2 different sessions, and the results are plotted in Figure 7c,d. The intersubject coefficient of variance were 22% (GM) and 9% (WM) ( $n = 5$ ) for amideCEST (Figure 7c). Excellent intrasubject agreement were found in the amideCEST measurements from 2 scans with intraclass correlation coefficient = 0.98 (GM) and 0.99 (WM) ( $n = 3$ ; Figure 7e).





**Figure 7: The differences in amideCEST signal between GM and WM**

GM and WM distributions of fitted amideCEST signal differences ( $\Delta Z_{amide}$ ) as well as the corresponding apparent relaxation rates (b:  $R_{amide}$ ) for all subjects ( $n = 5$ ). GM and WM regions show very similar distribution of  $\Delta Z_{amide}$ . After correction with the rotating frame relaxation method, WM regions show slightly higher contrast in the rotating frame relaxation rate  $R_{amide}$ . Comparison of fitted amideCEST (c:  $\Delta Z_{amide}$ ) and the corresponding rotating frame relaxation rates (d:  $R_{amide}$ ) in GM and WM for 5 healthy subjects. Comparison of amideCEST (e:  $\Delta Z_{amide}$ ) and the rotating frame relaxation rates (F:  $R_{amide}$ ) in GM and WM for the same subject ( $n = 3$ ; subject 2, 4 and 5) at 2 different sessions that is, scan1(S1) and scan 2 (S2) over 2 days. The error bar is the SD for each subject. GM, gray matter; WM, white matter

## Chapter 4 Discussion

We applied radial CEST acquisition, postprocessing, and quantification methods to isolate the amideCEST signal in human brain at 3T. We demonstrated that the amideCEST signal can be extracted with the updated PLOF method at 3T. The amideCEST images acquired with starCEST showed much improved image quality and SNR compared to the conventional 3D TSE CEST method with the same scan time. In order to achieve rapid acquisition for the TSE CEST, high TSE, and compressive sensing factors had to be used. The amideCEST maps acquired by the starCEST were robust to the brain movements (Figure 6), whereas motion corrections are required for the CEST maps with the TSE readout.

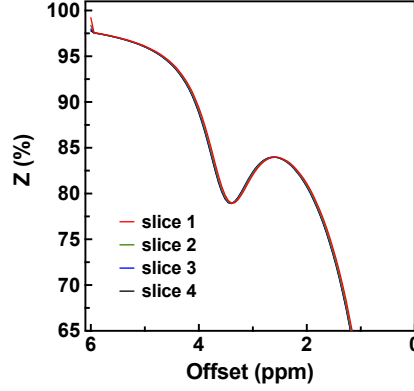
In contrast to a recent creatine validation study<sup>44,62</sup> using a knockout mouse model, validation of the amideCEST signal in the brain Z-spectrum cannot be easily achieved by removing the mobile proteins in tissue. However, we could exploit the property that CEST signals from amide and Guan protons are strongly pH-dependent, whereas rNOEs have weak pH dependence and MTC is insensitive to pH<sup>59,63</sup> in the physiological range. As demonstrated by the egg white model system (Figure 2), the amide CEST signal is visible at 3T with a peak around 3.5 ppm. However, there is still a large residual signal at 3.5 ppm (purple arrow in Figure 2b) that is not pH-dependent over the physiological range of pH 6.0-7.5 (Figure 2b). This conclusion is consistent with previous studies with animal stroke models at both high<sup>26,60,63</sup> and low MRI fields.<sup>64</sup> The source of the strong

residual background signal (purple arrow in Figure 2b) in addition to amideCEST was further investigated with hydrogen/deuterium exchange studies on egg white solutions. As seen from Figure 2c, a large portion of this broad background signal between 0 and 6 ppm disappeared for 70%  $D_2O$ , which we tentatively attribute to faster exchanging protons from amine and hydroxyl groups, signals of which are partially merged with water due to being in the intermediate to fast exchange regime at 3T. Due to this exchange regime, this portion of the CEST signal only has a weak pH dependence in the physiological range (6-7.5). In addition to this, the background (purple arrow in Figure 2b) may have contributions from rNOEs, for example, amide-rNOE and aromatic-rNOE. The existence of NOE components at 3.5 ppm has been suggested by hydrogen/ deuterium exchange studies,<sup>65,66</sup> which have long been used to probe protein structures. The exchange rates of the amide protons depend on solvent accessibility and hydrogen bonding. The amide protons that are part of an intramolecular hydrogen bond or buried from the water will exchange slowly or do not exchange at all, whereas an amide on the surface of protein and side chains will exchange rapidly with water. Hence, NOE signals at approximately 3.5 ppm may be attributable to non-exchanging amide protons or amide proton with extremely slow exchange rates, but they will still be able to transfer magnetization to water following a 2-step relayed NOE process with faster exchanging neighboring protons.<sup>67</sup> However, this hypothesis and the contribution from faster exchanging amine and OH groups to the pH-independent background for the 3T exchange regime still needs to be further validated and is beyond the scope of the current study. The  $D_2O$  study seems to indicate the contribution of the latter may be substantial. One recent study in which the myelin lipids in brain were removed also suggests that the rNOE signal from lipids may contribute to the positive side of the Z-spectrum.<sup>68</sup> From the egg white study, the background CEST signal at 3.5 ppm, extracted using either Lorentzian fitting<sup>8,34-37</sup> or asymmetry analysis,<sup>69,70</sup> includes many

components such as amineCEST, hydroxylCEST, and rNOE signals. All these effects, including both CEST and rNOE effects, but not exclusively, also apply to mobile proteins *in vivo*. Hence, the signal at 3.5 ppm extracted by the conventional Lorentzian fitting or asymmetry analysis can still be applied for the protein profiling in tissues such as tumor. When probing pH changes in tissue at the physiological range, the PLOF method can thus provide a quite sensitive contrast by extracting the major pH-dependent component at physiological pH, that is, amideCEST, or at higher  $B_1$ , amineCEST, as shown by Harris et al.<sup>71</sup>

In the starCEST method, the steady-state saturation is a consequence of both the saturation pulses and the MRI readouts during the mixing period. As previous steady-state CEST optimization demonstrated, higher flip angles of the MRI readout pulses reduce the CEST contrast, whereas lower flip angles lead to low SNR for the MRI images.<sup>72,73</sup> In the current study, a flip angle of 10 degrees was selected as a good compromise. The time to reach steady-state in starCEST when switching saturation offset is determined by the water rotating frame relaxation time, which is close to water  $T_1$  at low saturation powers.<sup>13</sup> As the simulations in Figure 8 show, the non-steady-state character of the saturation when switching irradiation frequency impacts mainly the initial offsets for slices 1-3. This can be easily solved by removing the first offset from the PLOF fitting.

Here, we demonstrated that, with same scan time, the starCEST approach can generate much better image quality than the 3D TSE CEST. The acquisition speed for the 3D MRI readout can be further improved with a 3D gradient- and spin-echo readout. Previously, a scan time of 20 s for each offset has been applied for the 3D gradient- and spin-echo CEST to achieve reasonable image quality.<sup>74</sup> This is much slower than the starCEST method and, similar to the TSE CEST approach, has the issue of not being robust to motion. Recently, several other whole brain CEST methods

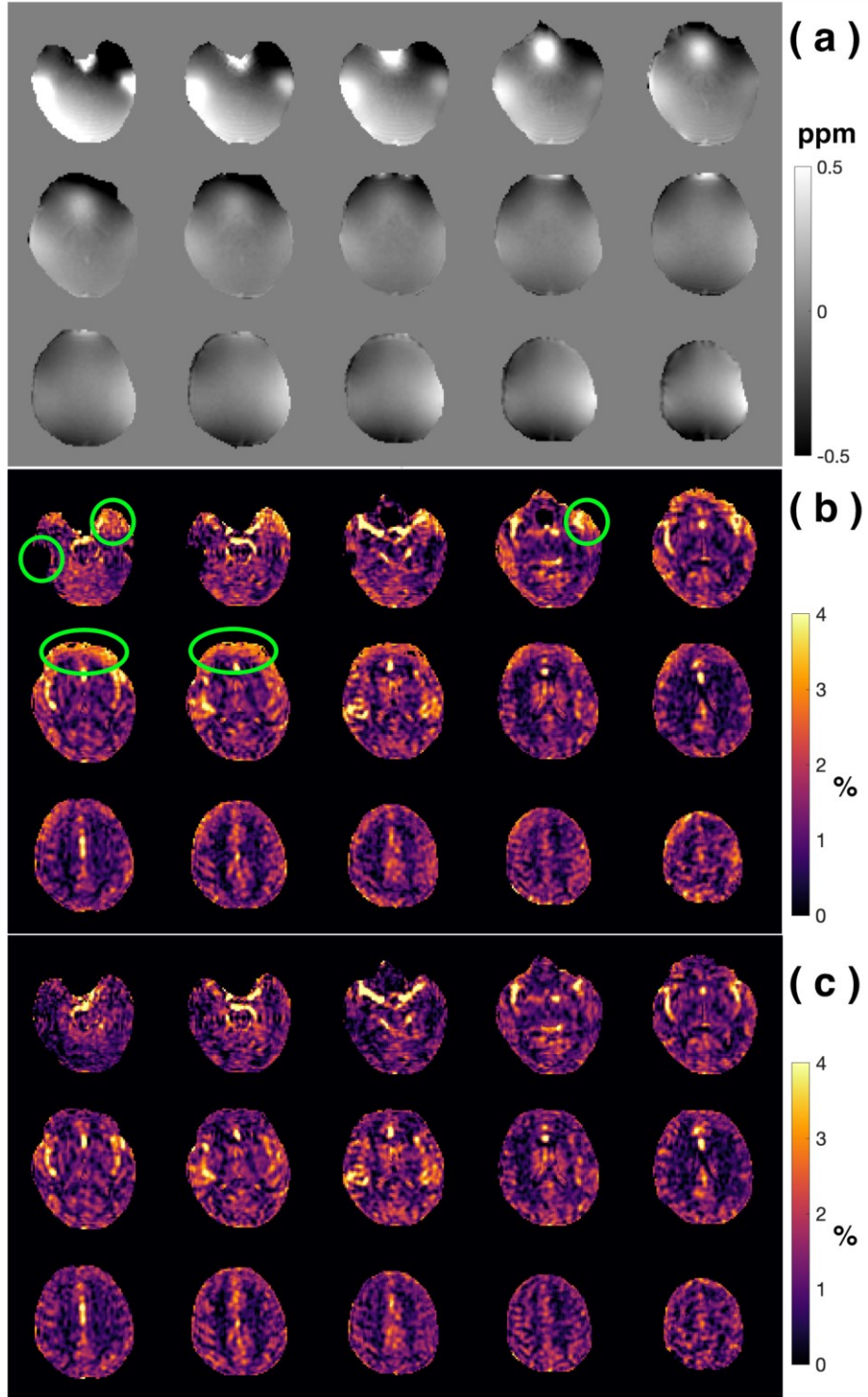


**Figure 8: Multi-slice simulation of the Z-spectrum**

Simulation of the Z-spectrum for each slice in the starCEST to demonstrate the impact of the non-steady-state character of the saturation. Only the first four slices are shown for clarity. In the simulation, the 50 ms since-Gauss saturation pulses and 53 ms mixing times were applied. The frequency was swept from 6 ppm to 1 ppm.  $T_1 = 1\text{ s}$ ,  $T_2 = 80\text{ ms}$  were set for the water pool, while  $T_1 = 1\text{ s}$ ,  $T_2 = 80\text{ ms}$ , exchange rate 160 Hz, offset 3.5 ppm and concentration 500 mM were assumed for the amide proton pool. The non-steady-state saturation significantly impacts the Z values in the slices 1-3 at the initial offset (6 ppm) and also shift the Z-spectrum of the first slice by 0.03 ppm.

have been developed to reach acquisition speeds comparable or faster than the starCEST method, including steady-state 3D EPI (11 s per offset),<sup>75</sup> snapCEST (6.9 s per offset),<sup>41,76</sup> and 3D EPI with controlled aliasing in parallel imaging results in higher acceleration (4.3 s per offset).<sup>37</sup> It is known that segmented 3D MRI readouts are very sensitive to motion, as suggested by arterial spin labeling MRI studies.<sup>77</sup> Thus, motion correction is still required for those 3D CEST MRI techniques, whereas starCEST is a multi-slice method and is inherently robust to motions. One advantage of the snap-CEST and 3D EPI CEST methods over the starCEST and steady-state 3D EPI methods is more flexibility in the CEST preparation module. Therefore, both snapCEST and 3D EPI CEST methods are more suitable for the CEST experiments with fast-exchanging protons.

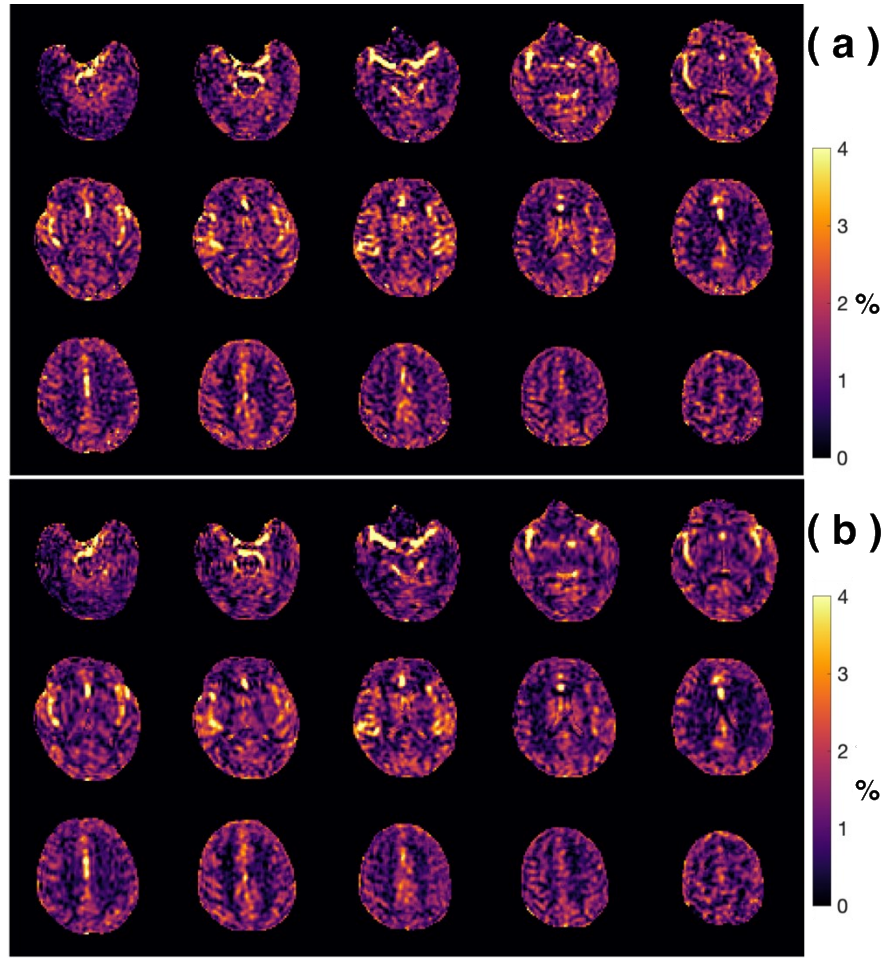
In the current study, the  $B_0$  maps were used for inhomogeneity correction in the amideCEST quantification. When the  $B_0$  shift is small ( $< 0.3\text{ ppm}$ ), the PLOF fitting method can find the amide peak and correct the  $B_0$  shift as demonstrated by Figure 9. It is still challenging for the PLOF method to correct the offsets with higher  $B_0$  shift values ( $> 0.3\text{ ppm}$ ). The pulsed CEST method is



**Figure 9: The effect of  $B_0$  correction in amideCEST maps**

(a)  $B_0$  map recorded with a dual-echo sequence. Comparison between the amideCEST maps extracted using the PLOF method without (b) and with (c)  $B_0$  correction. Hyper/hypo intensity regions in the amideCEST maps are indicated with green circles and are due to the high  $B_0$  shifts (> 0.3 ppm).

robust to  $B_1$  inhomogeneity compared with continuous wave CEST because 20% in the peak  $B_1$  value change is only equivalent to an 10.8% change in continuous wave CEST ( $0.54 \times 20\%$ ). As the amideCEST maps closely resemble to the maps obtained by suppressing MTC with variable delay multi-pulsed method<sup>79</sup> in which strong GM signal was observed in the amide proton transfer map. Although a clear difference in GM and WM was observed in the  $\Delta Z_{amide}$  maps (Figure 5c) as well as the Z-spectra in Figure 3d,e, the GM/WM contrast was highly variable between different subjects (Figure 7c). The intrasubject reproducibility study (Figure 7e) showed consistent GM/WM contrast. Further large-scale studies will be necessary to reveal the origin of the highly variable GM/WM contrast among different subjects. The amideCEST signal extracted by the PLOF method ( $< 2\%$ ) was much lower than the total signal at 3.5 ppm extracted with the Lorentzian fitting method.<sup>38,42,43</sup> Hence, the MLSVD method was used to enhance the overall SNR, as illustrated by the comparison study in Figure 10. Due to the low amideCEST signal, the background fitting in the PLOF method is critical for extracting reliable amideCEST values. Use of a third-degree polynomial function to fit the brain Z-spectrum at 3T was challenging due to the strong water DS as demonstrated in Figure 3d. As a consequence, a reduced fitting range ( $< 5.1$  ppm) had to be used for the PLOF method with the polynomial function, which underestimated the amideCEST values. The current fitting function (Equation 5) provided a much-improved background fitting over a wider frequency range (6.5 ppm) (Figure 3d), which lead to reliable amideCEST values as demonstrated by the intrasubject reproducibility study (Figure 7e). The proposed starCEST method not only works for the brain but is also expected to be well suited for study on body tissues, such as the liver and kidney, where motion is a serious issue as seen in our recent liver study.<sup>15</sup>



**Figure 10: The effect of MLSVD processing in amideCEST maps**

Comparison of the amideCEST maps without and with MLSVD post-processing. (a) The original amide CEST maps extracted using PLOF without MLSVD denoising. (b) MLSVD denoising applied with truncation numbers 48, 48 and 10.



## Chapter 5 Conclusion

We have demonstrated that a multi-slice starCEST method can be used for the imaging of amideCEST effect at clinical field strength. The PLOF method was implemented for quantifying the pH-sensitive amideCEST signal. The MLSVD post-processing method was applied to achieve enough SNR for the amideCEST maps, which trended toward slightly higher contrast in GM. This trend disappeared when calculating MTC- and T<sub>2</sub>-corrected rotating frame relaxation rates.

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